

Fourier Transform and Fourier Encoding in Magnetic Resonance

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Magnetic Field and Matter

empty space

$B = \mu_0 H$

Matter

$B = \mu_0 (H + M)$

Magnetic Susceptibility

$$B = \mu_0 (H + M) = \mu_0 (I + \chi) H$$

$$\chi = \chi_{diam.} + \chi_{param.} + \chi_{nuc. param.}$$

All organic substances. Effect of the orbital electrons within the magnetic field

$\chi_{H_2O} = -9.06 \cdot 10^{-6}$

Diamagnetic materials have $\chi_m < 0$, i.e., they are less permeable than free space to magnetic fields.

Paramagnetic materials have $\chi_m > 0$, i.e., they are more permeable than free space to magnetic fields.

Unpaired electrons have a magnetic moment (μ_B)

Magnetic moment of atoms with an odd number of nucleons (μ_K)

$\mu_K = \mu_B / 1836$

BioImaging, MR_gf (MUG 2004)

Spins in the Atom (Nucleus)

If a charged particle has non-zero angular momentum I , then it also has a magnetic moment μ (and vice versa), and $\mu \parallel I$. I depend on the spin quantum number s which again depend on the mass number A .
 Odd A : s may be $1/2$ (e.g. protons), $3/2$, $5/2$, ... Even A : s may be 0 , 1 , 2 , ...

$$|I| = \sqrt{s(s+1)}\hbar$$

e.g. Proton

„gyroscope“

The ratio μ/I is called the *magnetogyric ratio* γ and is specific for different kinds of nuclei

$$\gamma = \mu / I \qquad \hbar = \frac{h}{2\pi}$$

BioImaging, MR_gf (MUG 2004)

$B=0$ Random Orientation

$T > 0$

Orientation in Magnetic Field

$T > 0$

$\omega_0 = \gamma B_0$

Macroscopic Magnetization in Equilibrium (M_0)

from Boltzmann statistics:

$$M_0 = \frac{\rho \gamma^2 \hbar^2 B_0}{3kT}$$

$\omega_0 = \gamma B_0$

Excitation

Transversal component of an RF-magnetic field precesses synchronic (ω_0) with the nuclear spins

constant torque on macroscopic magnetization M_0

$\omega_0 = \gamma B_0$

linear polarized Field

Magnetization M will be flipped by the constant torque produced by B_1

$\alpha = \gamma B_1 t_p$

$\omega_0 = \gamma B_0$

NMR -Signal

$FID \propto M_{xy}$

induced voltage (FID)

Relaxation

HF 90°

$M_z = M_0(1 - e^{-t/T_1})$

$M_{xy} = M_0 e^{-t/T_2^*}$

Chemical Shift & FID

The resonance frequency $\omega(\mathbf{r})$ is proportional to the magnetic flux density $\mathbf{B}_{mac}(\mathbf{r})$ at the position of the nucleus:

$$\mathbf{B}_{mac}(\mathbf{r}) \propto \mathbf{B}_{mac}(\mathbf{r}) - \sigma \mathbf{B}_0$$

\mathbf{B}_{mac} is the macroscopic flux density which results after an object with magnetic susceptibility χ_i is placed into the main field \mathbf{B}_0 . \mathbf{B}_{mac} depends on \mathbf{B}_0 , on the geometry of the object and on the susceptibility distribution internal χ_i and external χ_e of the object. The electrons of different molecules (chemical environment) "shield" the nucleus to varying degrees σ depending on the position of the nucleus in the molecule. This is the chemical shift effect.

FID

interference of different frequencies

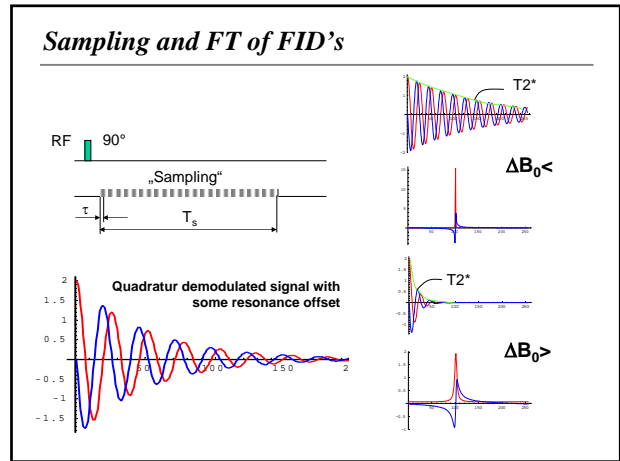
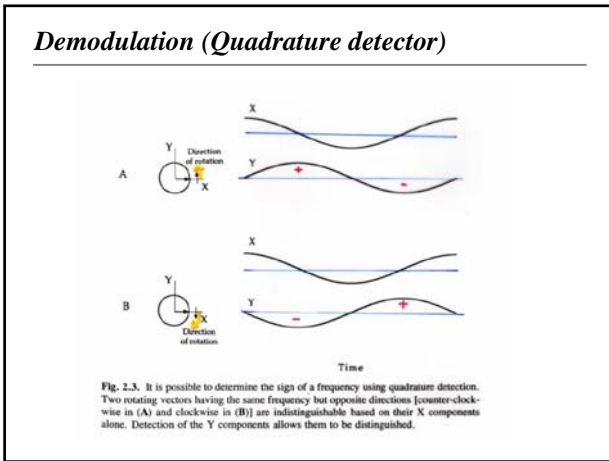
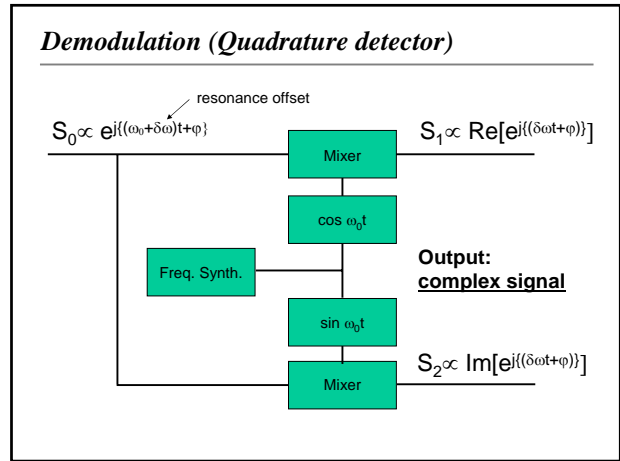
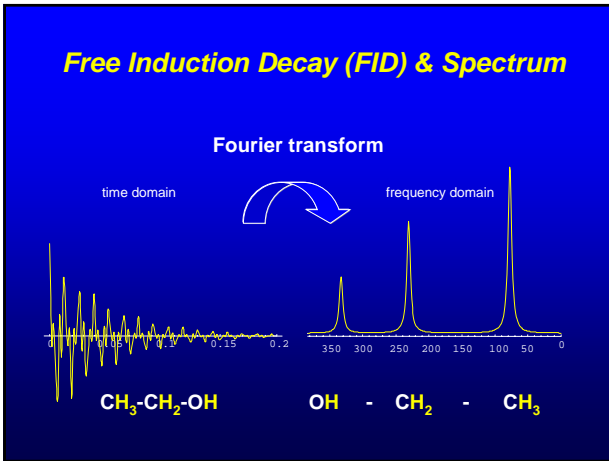
BioImaging, MR, gi (MUG 2004)

Chemical Shift Within a Molecule

$\text{CH}_3\text{-CH}_2\text{-OH}$ f_{CH_3}

$\text{CH}_3\text{-CH}_2\text{-OH}$ f_{CH_2}

$\text{CH}_3\text{-CH}_2\text{-OH}$ f_{OH}



Fourier Transform (FT)

- The one-dimensional Fourier Transform is given by

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt = \int_{-\infty}^{\infty} f(t) (\cos \omega t - i \sin \omega t) dt$$
- $F(\omega)$ is in general a complex number, with real and imaginary parts

$$F(\omega) = R(\omega) + iI(\omega)$$

$$= M(\omega) e^{i\phi(\omega)}, \quad \text{Magnitude } M(\omega) = \sqrt{R^2(\omega) + I^2(\omega)}$$

$$\text{Phase } \phi(\omega) = \tan^{-1} \left[\frac{I(\omega)}{R(\omega)} \right]$$

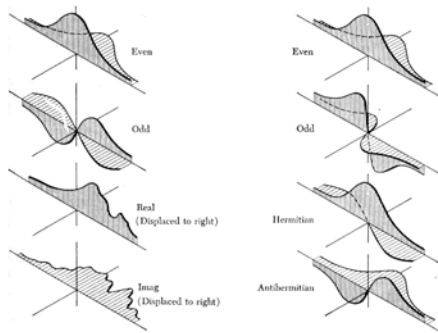
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Some Properties of Fourier Transform

- Linearity
 - $\mathcal{F}[af(t) + bg(t)] = a \mathcal{F}[f(t)] + b \mathcal{F}[g(t)]$
- Scaling
 - $\mathcal{F}[f(at)] = \frac{1}{|a|} \mathcal{F}[f(t/a)]$
- Shifting
 - $\mathcal{F}[f(t-a)] = e^{-i\omega a} \mathcal{F}[f(t)]$
- Convolution theorem
 - $\mathcal{F}[f(t)g(t)] = \mathcal{F}[f(t)] * \mathcal{F}[g(t)]$ resp. $\mathcal{F}[f(t) * g(t)] = \mathcal{F}[f(t)] \mathcal{F}[g(t)]$
- Parseval's theorem (Energy content)
 - $\int_{-\infty}^{\infty} |f(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(\omega)|^2 d\omega$

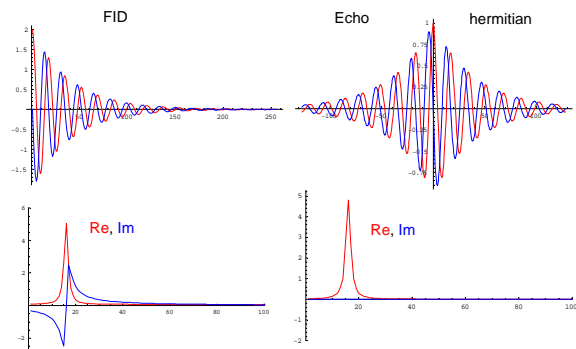
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Some Properties of Fourier Transform



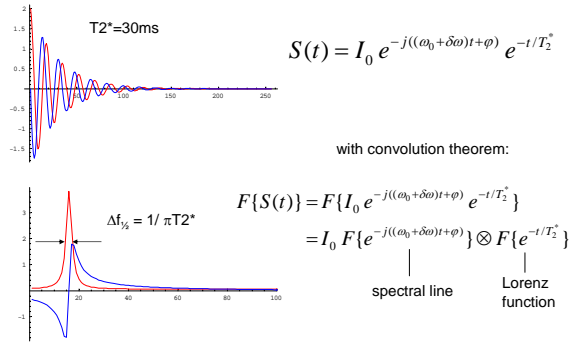
Biolmaging, MR, gf (MUG 2004)

Some Properties of FT in NMR



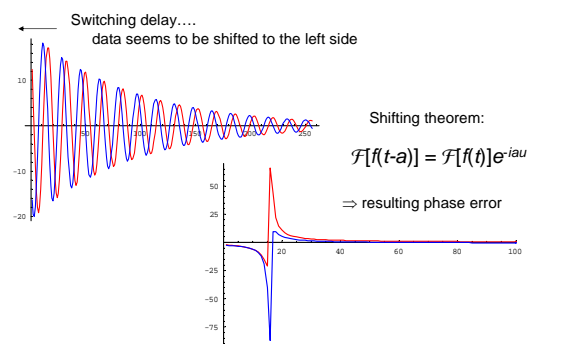
Biolmaging, MR, gf (MUG 2004)

Some Properties of FT in NMR

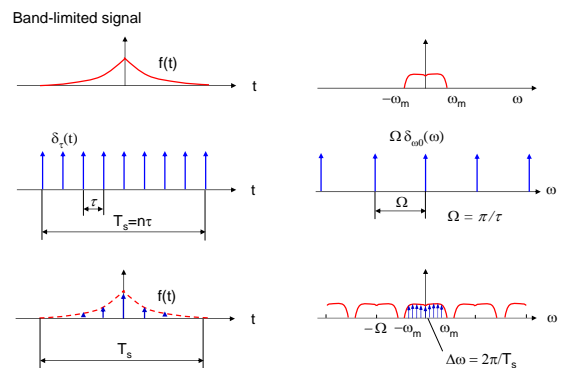


Biolmaging, MR, gf (MUG 2004)

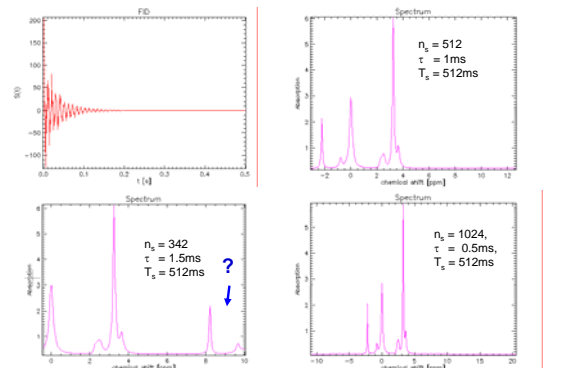
Some Properties of FT in NMR



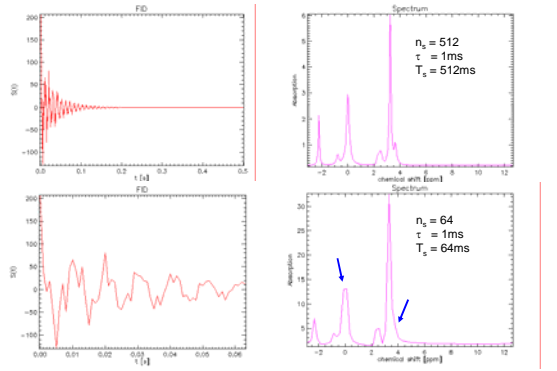
Basic Properties of DFT



DFT in MRS: Bandwidth = $2\pi/\tau$



DFT in MRS: Spectral Resolution = $2\pi/T_s$

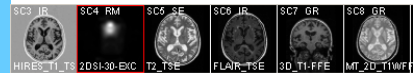


What do we measure in a MR-imaging experiment?



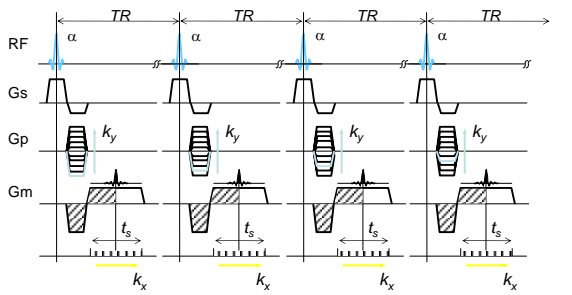
Patient: different NMR parameters (e.g. T1, T2, δ ...) depending on type of tissue: white/gray matter, tumor...
 ...tissue composition: blood volume, fat content...
 ...state of tissue: inflammation, activation, oxygenation, temperature...
 ...motion (blood flow, heart beat, ...)

Pulse-Sequences:
 ⇒ encode the image
 ⇒ define the contrast



Pulse - Sequences

Pulse-Sequences: ⇒ encode the image ⇒ define the contrast



Bioimaging, MR, gr (MUG 2004)

Signal Localization: Slice Selection

$$\omega_0 = \gamma B_0$$

$$\text{Excitation: } \omega_{RF} = \omega_0$$

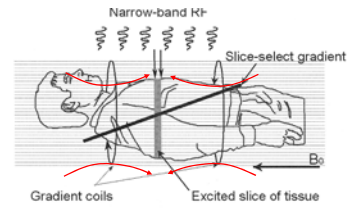
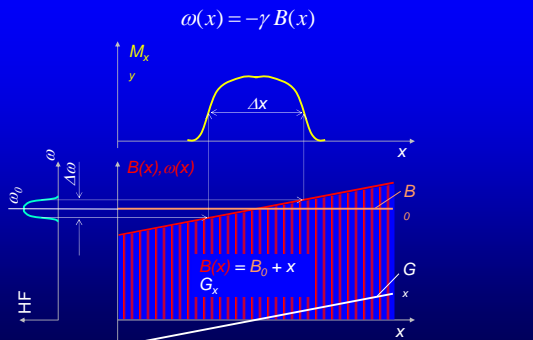


FIGURE 15-4. The slice select gradient (SSG) disperses the precessional frequencies of the protons in a known way along the gradient. A narrow-band radiofrequency (RF) pulse excites only a selected volume (slice) of tissues, determined by RF bandwidth and SSG strength.

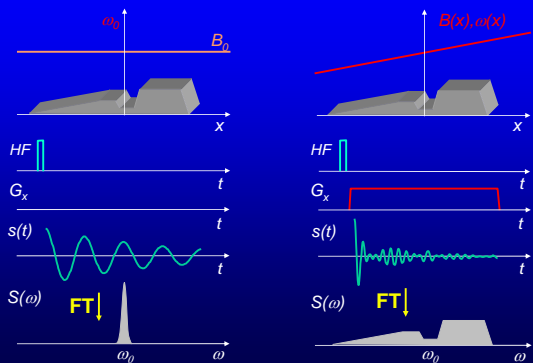
Bioimaging, MR, MRI (MUG 2004)

Slice Profile \propto Power Spectrum RF



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Frequency encoding $\omega(x) = -\gamma B(x)$



Bioimaging, MR, gr (MUG 2004)

Spin Wrap Imaging: Signal Encoding

(2d Fourier Method)

$$\phi_x = -\gamma x \int_0^t G_x d\tau \equiv k_x x \quad \phi_y = -\gamma y \int_0^t G_y d\tau \equiv k_y y$$

Biomed Imaging, MR, gl (MUG 2004)

Signal & k-space

$$s(t, G_y) = \iint I(x, y) e^{-j[\phi(x,t) + \phi(y, G_y)]} dx dy$$

$$\phi(x, t) = -\gamma x \int_0^t G_x d\tau \equiv k_x x$$

$$\phi(y, G_y) = -\gamma y \int_0^t G_y d\tau \equiv k_y y$$

$$s(k_x, k_y) = \iint I(x, y) e^{-j(k_x x + k_y y)} dx dy$$

...Signal equation = Fourier integral

$$\tilde{I}(x, y) = \frac{1}{2\pi} \iint s(k_x, k_y) e^{j(k_x x + k_y y)} dk_x dk_y$$

...discrete data: 2DFT

$$\hat{I}(p, q) = \frac{1}{NM} \sum_{n=-\frac{N}{2}}^{\frac{N}{2}-1} \sum_{m=-\frac{M}{2}}^{\frac{M}{2}-1} s(n, m) e^{j2\pi(\frac{n}{N}p + \frac{m}{M}q)}$$

k_x, k_y encoded signals

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Measured data "k-space" 2D-FFT Image I(x,y)

Each point within the k-space contribute to the image

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Image reconstruction

Center of k-space determines signal intensity & main image contrast

Biomed Imaging, MR, gl (MUG 2004)

Image reconstruction

out parts of k-space determines the resolution

Biomed Imaging, MR, gl (MUG 2004)

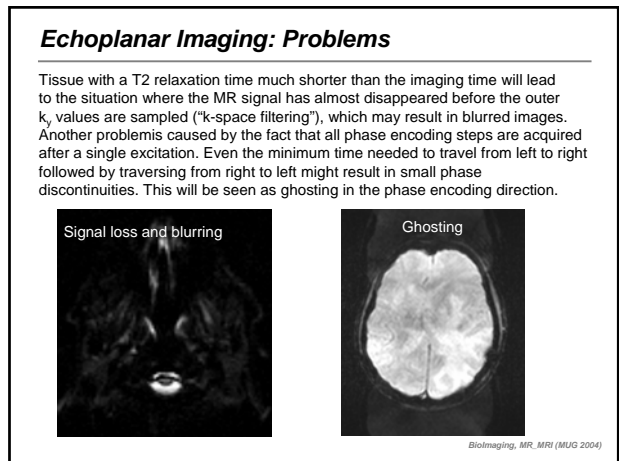
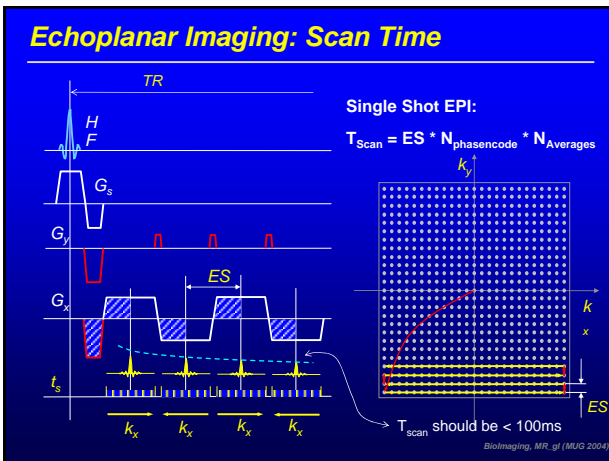
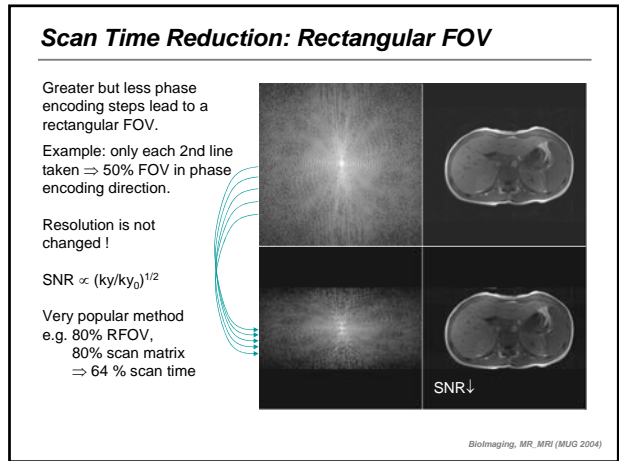
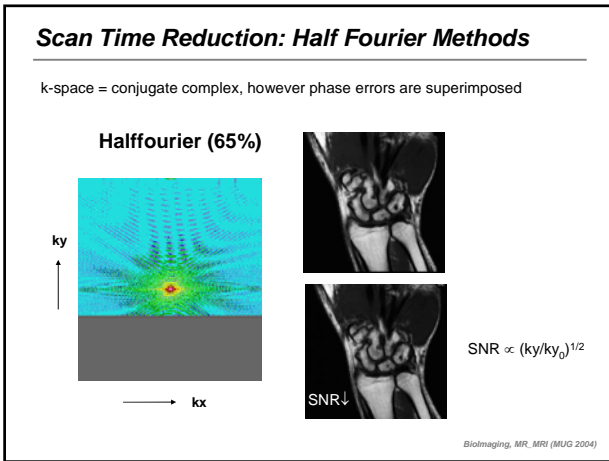
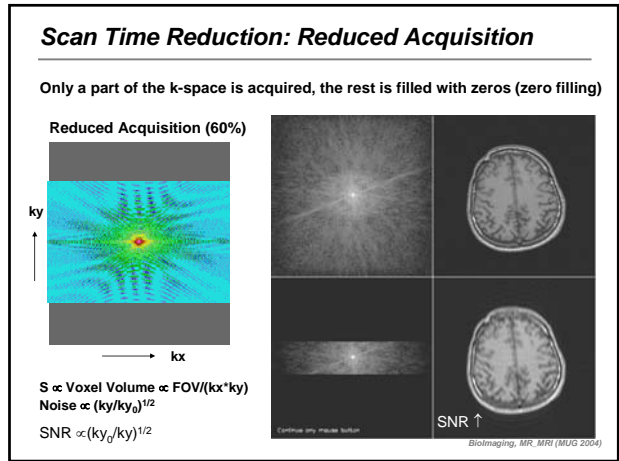
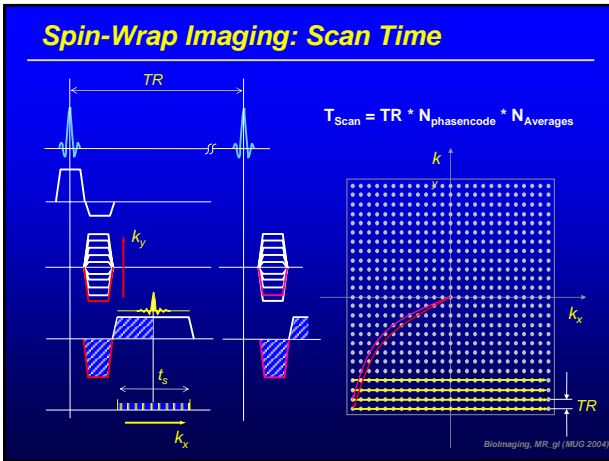
Point spread function in MRI

Point spread function answers the question, how much blurring would occur if you are trying to image a point.

Data always contains the effect of truncation and sampling

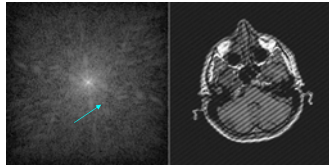
$$f(t) = f(t) \text{rect}(t/T_s) \quad F(\omega) = F(\Delta\omega) \otimes T_s \frac{\sin(2\pi T_s \omega)}{2\pi T_s \omega}$$

Biomed Imaging, MR, gl (MUG 2004)

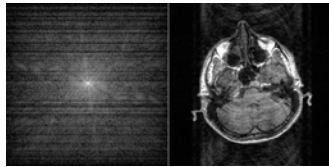


Artifacts Fourier MRI

Single spike in raw data
(e.g. spark due to low humidity)



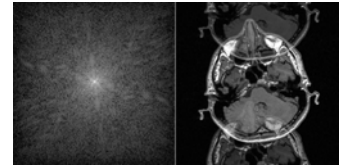
Random amplitude error



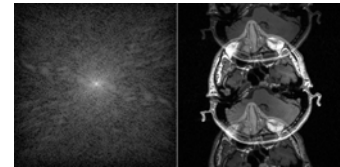
Biomedaging, MR, MRI (MUG 2004)

Artifacts Fourier MRI (N/2 Ghost)

Phase error each 2 line
($\pi/3$)



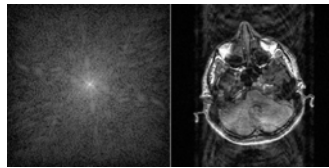
Shift each 2 line -0.5 samples



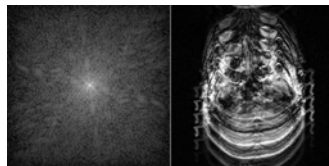
Biomedaging, MR, MRI (MUG 2004)

Artifacts Fourier MRI

random phase error
(motion)



periodic phase error
(e.g. periodic physiologic motion)



Biomedaging, MR, MRI (MUG 2004)